Physics 5645

Quantum Mechanics A

Problem Set IX

Due: Thursday, Nov 29, 2018 (DEADLINE EXTEDNED TO TUESDAY, DEC. 4)

- 9.1 Constructing spherical harmonics.
- (a) Use the fact that $L_z|11\rangle = \hbar|11\rangle$ and $L_+|11\rangle = 0$ and the position representations of L_+ and L_z ,

$$\langle \vec{r} | L_{+} | \psi \rangle = \hbar e^{i\phi} \left(\frac{\partial}{\partial \theta} + i \cot \theta \frac{\partial}{\partial \phi} \right) \langle \vec{r} | \psi \rangle,$$
$$\langle \vec{r} | L_{z} | \psi \rangle = \frac{\hbar}{i} \frac{\partial}{\partial \phi} \langle \vec{r} | \psi \rangle,$$

to explicitly derive $Y_1^1(\theta, \phi)$. Normalize it by carrying out the appropriate spherical integral. To be consistent with the usual convention be sure to include the appropriate factor of $(-1)^l$.

(b) By repeatedly applying the lowering operator L_{-} using the property that $L_{-}|l,m\rangle = \hbar\sqrt{(l+m)(l-m+1)}|l,m-1\rangle$ and the position representation of L_{-} ,

$$\langle \vec{r} | L_{-} | \psi \rangle = -\hbar e^{-i\phi} \left(\frac{\partial}{\partial \theta} - i \cot \theta \frac{\partial}{\partial \phi} \right) \langle \vec{r} | \psi \rangle,$$

obtain $Y_1^0(\theta,\phi)$ and $Y_1^{-1}(\theta,\phi)$. Compare your results with the familiar expressions,

$$Y_1^0(\theta,\phi) = \left(\frac{3}{4\pi}\right)^{1/2} \cos\theta, \quad Y_1^{\pm 1}(\theta,\phi) = \mp \left(\frac{3}{8\pi}\right)^{1/2} \sin\theta e^{\pm i\phi}.$$

9.2 Parity and angular momentum.

Prove that under the parity operation $\vec{r} \rightarrow -\vec{r}$,

$$Y_l^m \to (-1)^l Y_l^m$$
.

Hint: Show that this is true for Y_l^l , for which you have a simple explicit form, and then verify that applying L_- does not alter the parity.

9.3 Spin-1 rotation matrix.

Recall that the matrix representation of J_x for a spin-1 particle is,

$$J_x^{(1)} = rac{\hbar}{\sqrt{2}} \left(egin{array}{ccc} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{array}
ight).$$

- (a) Compute the matrices $J_x^{(1)^2}$ and $J_x^{(1)^3}$. You should find that $J_x^{(1)^3} = \hbar^2 J_x^{(1)}$.
- (b) By Taylor expanding the exponential, and using the results of Part (a), show that the spin-1 matrix representation of the rotation operator for an x-axis rotation by angle ϕ , $D(R(\phi \hat{i})) = e^{-i\frac{J_x}{\hbar}\phi}$, is

$$e^{-i\frac{J_x^{(1)}}{\hbar}\phi} = 1 - i\sin\phi\left(\frac{J_x^{(1)}}{\hbar}\right) + (\cos\phi - 1)\left(\frac{J_x^{(1)}}{\hbar}\right)^2.$$

- 9.4 Spherical harmonics in Cartesian coordinates.
- (a) Show that the l=1 spherical harmonics (see 9.1) can be expressed in Cartesian coordinates as,

$$Y_1^0 = \left(\frac{3}{4\pi}\right)^{1/2} \frac{z}{r}, \quad Y_1^{\pm 1} = \mp \left(\frac{3}{4\pi}\right)^{1/2} \frac{x \pm iy}{2^{1/2}r}.$$

(b) Now consider a particle in a state for which the position-space wave function is,

$$\psi(\vec{r}) = A(x + y + 2iz)e^{-\alpha r},$$

where A is a normalization constant. Using the result of Part (a), determine the probabilities for all possible results of an L_z measurement on this particle.

9.5 Rotating Y_l^m s.

Under a rotation through angle ϕ about the x-axis the x, y, and z components of the position vector of a particle transform as

$$x \to x$$

 $y \to y \cos \phi - z \sin \phi$
 $z \to z \cos \phi + y \sin \phi$

2

Thus, under this rotation, the spherical harmonic Y_1^0 must transform as

$$Y_1^0 = \left(\frac{3}{4\pi}\right)^{1/2} \frac{z}{r} \to \left(\frac{3}{4\pi}\right)^{1/2} \frac{z\cos\phi - y\sin\phi}{r}$$

(The difference in sign of the $y \sin \phi$ term is due to the fact that we are considering an *active* rotation.)

- (a) Expand the rotated Y_1^0 in terms of the unrotated Y_1^1 , Y_1^0 , and Y_1^{-1} .
- (b) Compare your result with that obtained using the spin-1 matrix representation of the x-axis rotation operator found in 9.3. **Hint:** Apply the rotation matrix to the column vector

$$\begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix},$$

which corresponds to the state Y_1^0 .

- (c) Repeat Parts (a) and (b) for the rotated Y_1^1 and Y_1^{-1} .
- 9.6 Spin-3/2 matrix representation of J_x .
- (a) Using

$$J_{\pm}|j,m\rangle = \hbar\sqrt{(j\mp m)(j\pm m+1)}|j,m\pm 1\rangle$$
, where $J_{\pm} = J_x \pm iJ_y$,

determine $J_x^{(3/2)}$, the 4×4 matrix representation of J_x for j=3/2 in the J_z basis.

(b) Find the eigenvalues of $J_x^{(3/2)}$ and confirm that they are $-\frac{3}{2}\hbar, -\frac{1}{2}\hbar, \frac{1}{2}\hbar, \frac{3}{2}\hbar$.